

Structure, Dynamics, and Phase Behavior of Lecithin-Based Microemulsions

Gerardo Palazzo



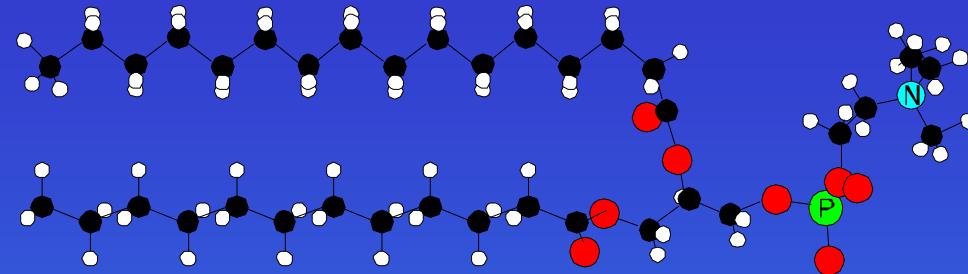
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ASSOCIATIONS IN SOLUTION: FUNCTION, PERFORMANCE, AND SYNTHESIS
(July 22-26 2007 Barga, Italy)

Lecithin = 1,2-diacyl-*sn*-glycero-3-phosphocholine



LS & SANS investigations*

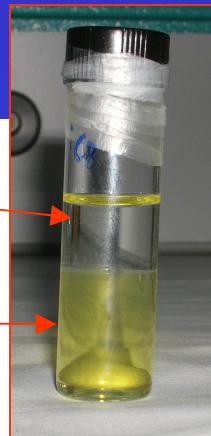
water addition induces the formation of giant cylindrical reverse micelles

Organogel: transparent, isotropic, highly viscous

* Schurtenberger et al. *J. Phys. Chem.* **1990**, *94*, 3695
Schurtenberger & Cavaco *Langmuir* **1994**, *10*, 100.
Schurtenberger et al. *Langmuir* **1996**, *12*, 2433.



pure oil



gel

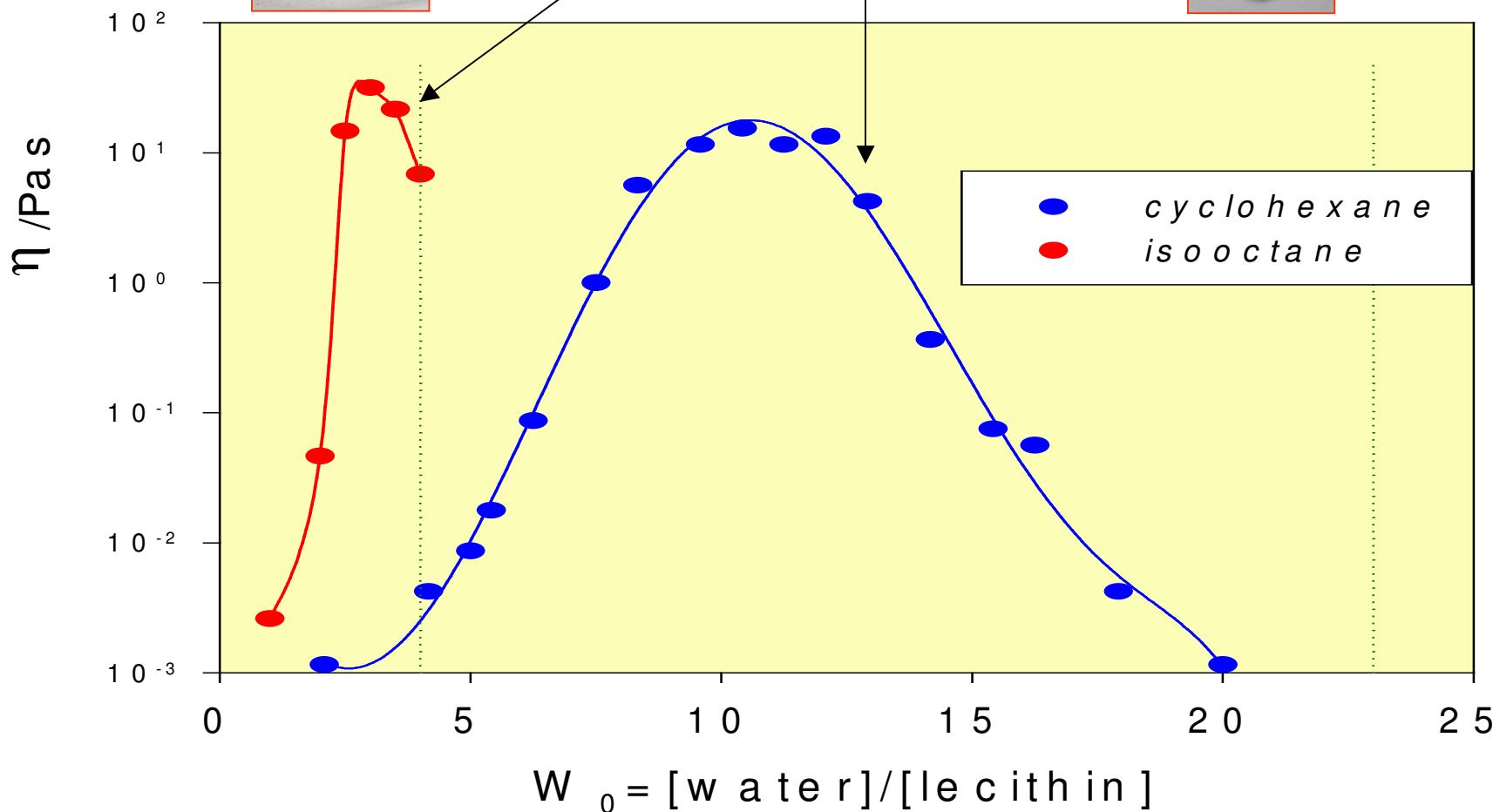
- different phase separation upon water dilution: Winsor II (C_6) or liquid-gas (iC_8)

- drop in viscosity

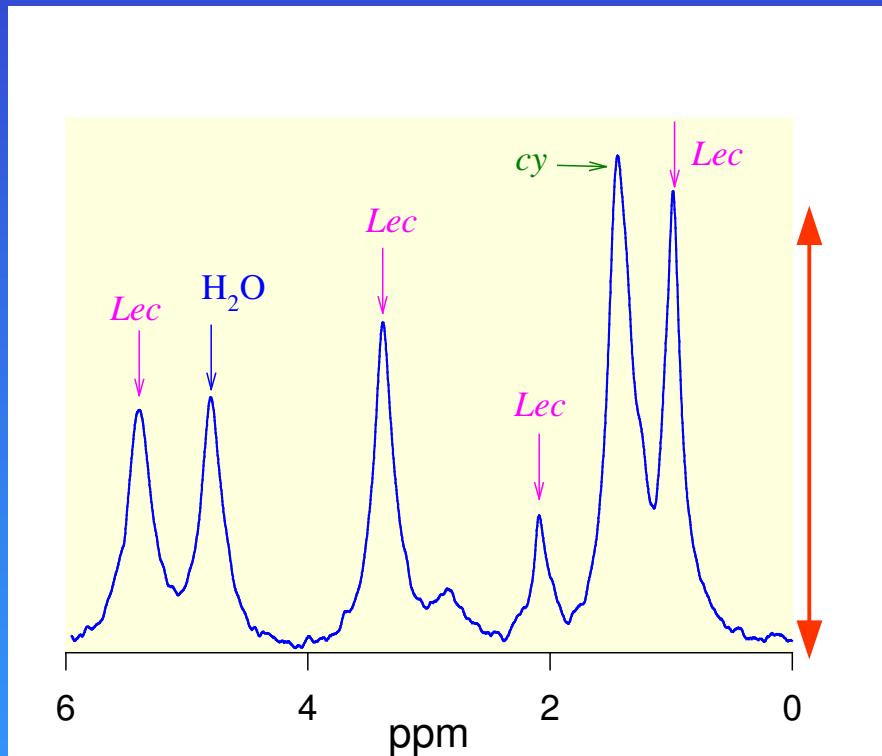
$$\Phi = 0.036$$

reverse micelles

water

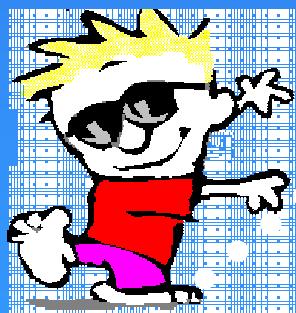


The Measurement of Molecular Motion Using PGSE-NMR



Echo attenuation $E(q,t)$

q & t : instrumental parameters



$$E(q,t) = \int P(Z,t) e^{iqZ} dZ$$

P(Z,t) Gaussian

Simple liquid

$$P(Z,t) = \frac{\exp\left(-\frac{Z^2}{4Dt}\right)}{\sqrt{4\pi Dt}}$$

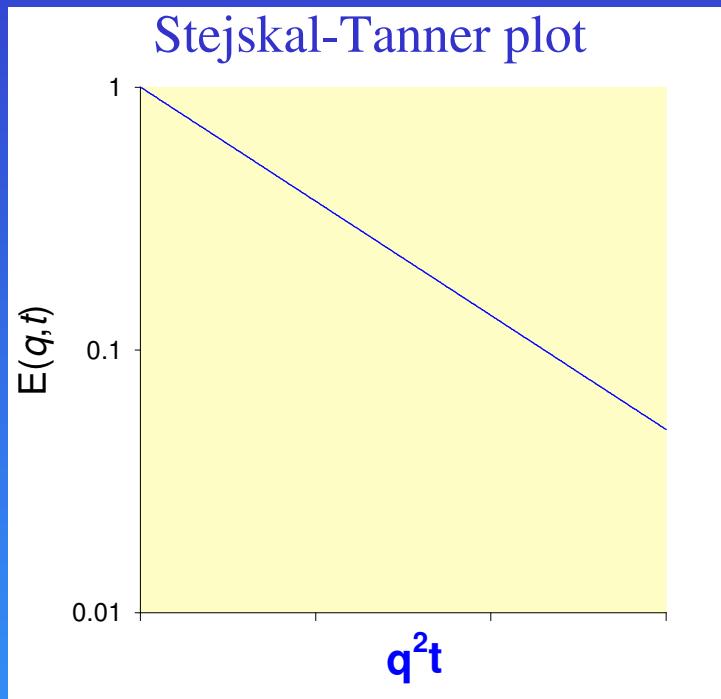


$$E(q,t) = \exp(-\mathbf{D}tq^2)$$

Unknown P(Z,t)

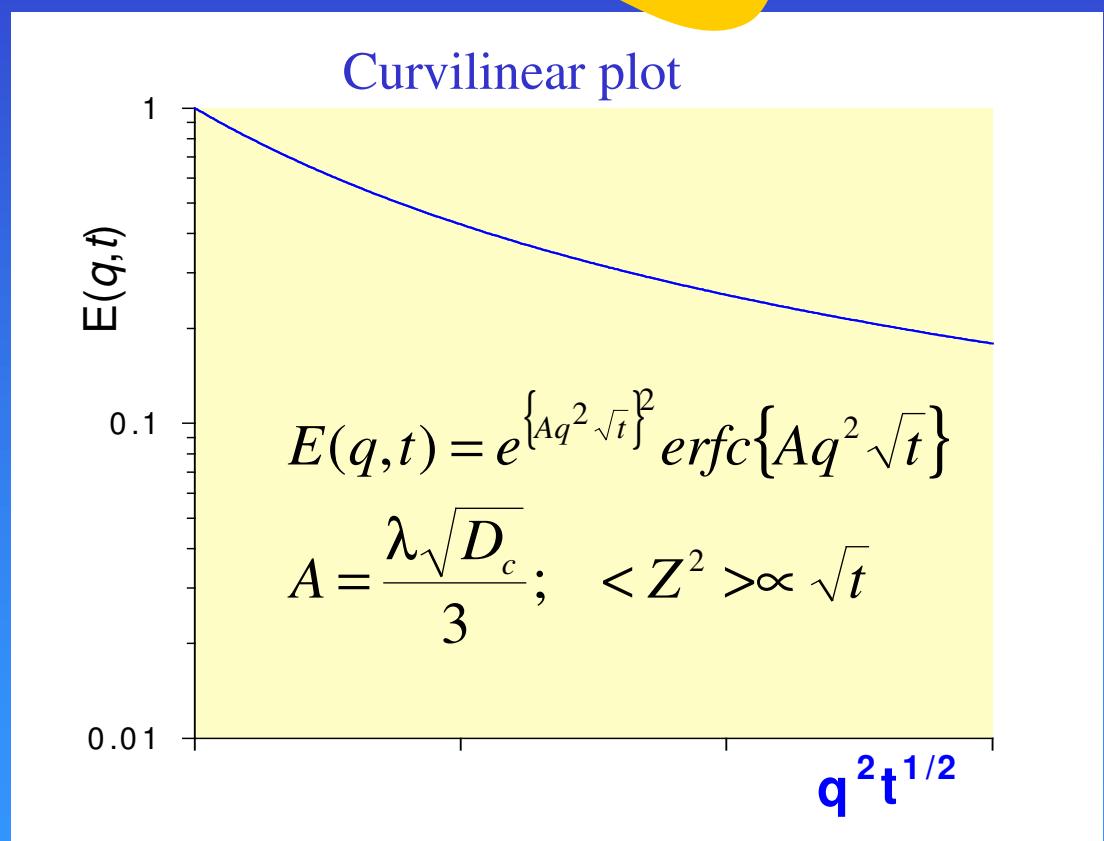
$$\lim_{q^2 \rightarrow 0} E(q,t) = 1 - \frac{\langle Z^2 \rangle q^2}{2}$$

Gaussian diffusion: Random Walk



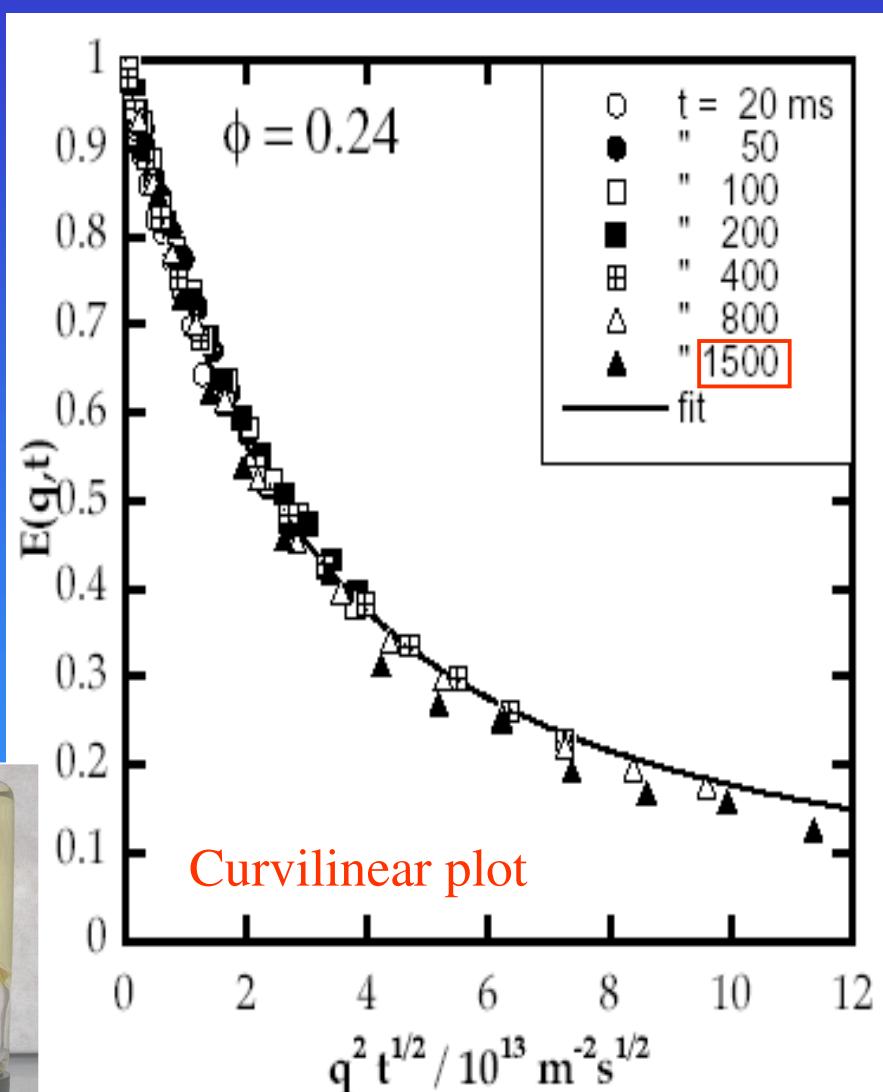
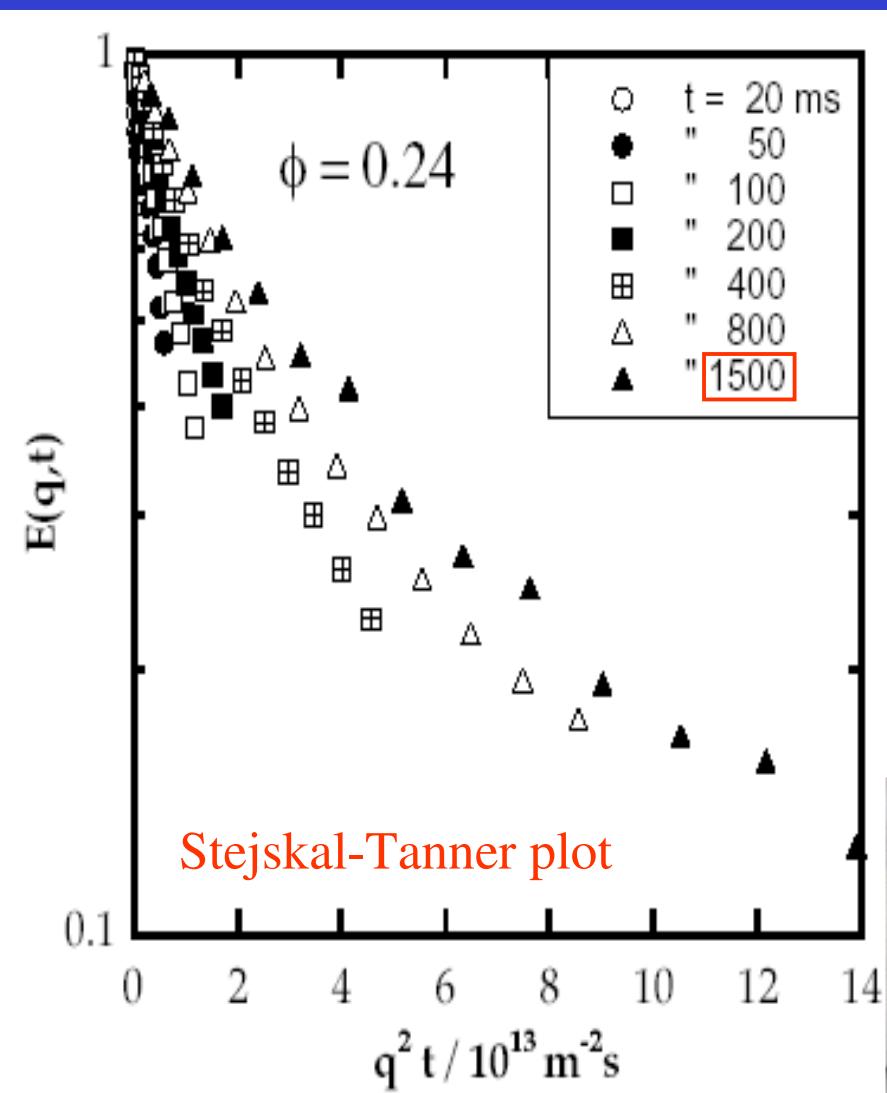
$$E(q,t) = e^{-Dq^2 t}$$
$$\langle Z^2 \rangle = 2Dt$$

Curvilinear diffusion: Random Walk on a Random Walk



Cyclohexane $W_0=10$

Angelico, Olsson, Palazzo, Ceglie *Phys. Rev. Letters* 1998, 81, 2823



In *cC6* lecithin experiences pure curvilinear diffusion
up to an observation time of 1.5 s

very long micelles (in 1.5 s lecithin explore about 2 μm)
micellar lifetime and the lecithin residence time > 1.5 s
disconnected micelles (not branched)

diffusion experiment in a branched network ?

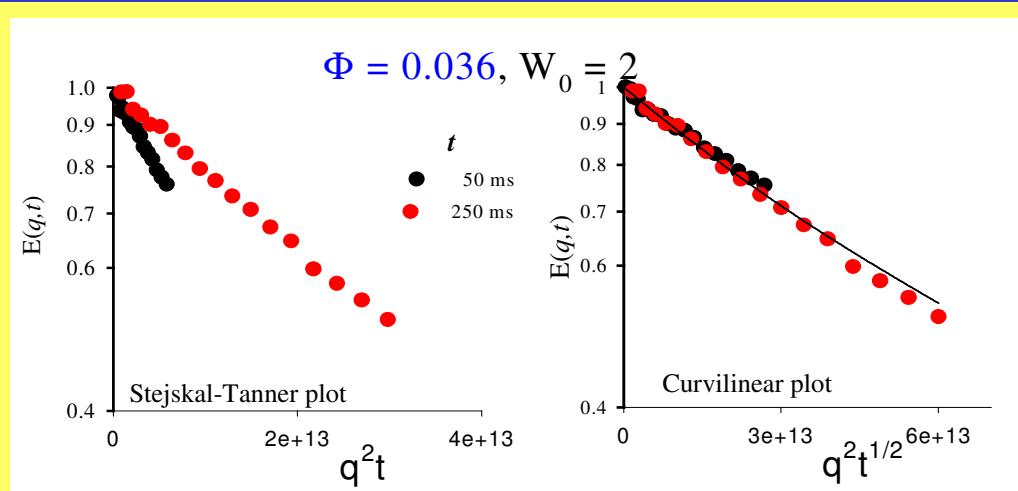
Curvilinear
 $\langle Z^2 \rangle \propto \sqrt{t}$

Random walk (Gaussian)
 $\langle Z^2 \rangle \propto t$



Ambrosone, Angelico, Ceglie, Olsson, Palazzo
Langmuir 2001, 22, 6822

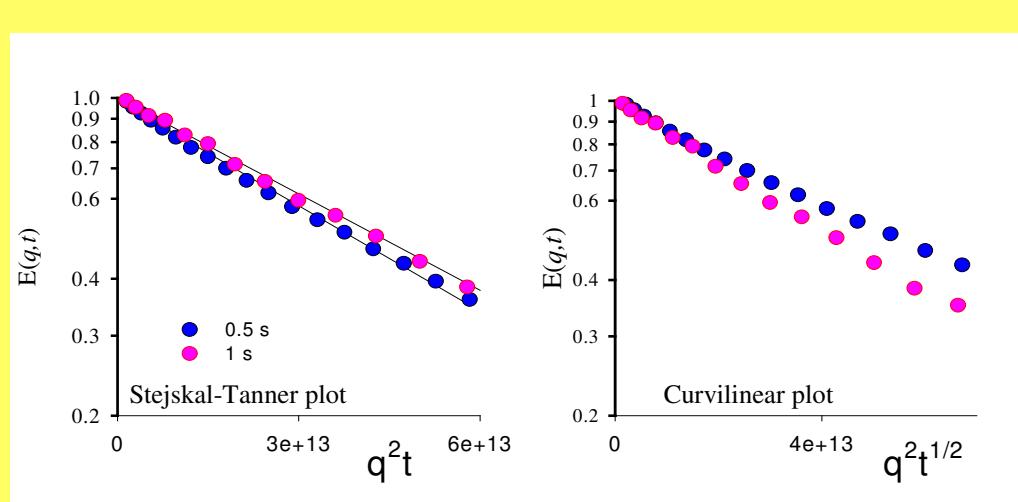
$$D = \frac{D_c}{3N}$$
$$N = \frac{L_{branch}}{2\lambda}$$



Low Φ , short times

curvilinear diffusion
 $\lambda = 150 \text{ \AA}$

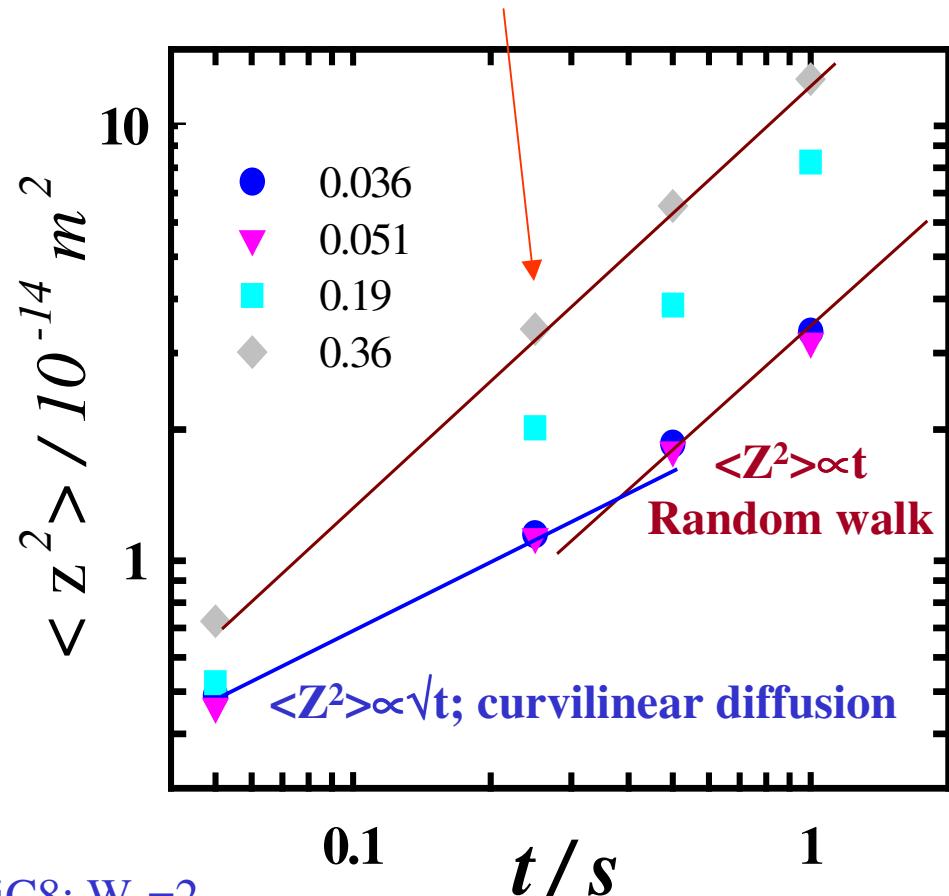
Isooctane, $W_0=2$ low concentration



Low Φ , long times

almost Gaussian

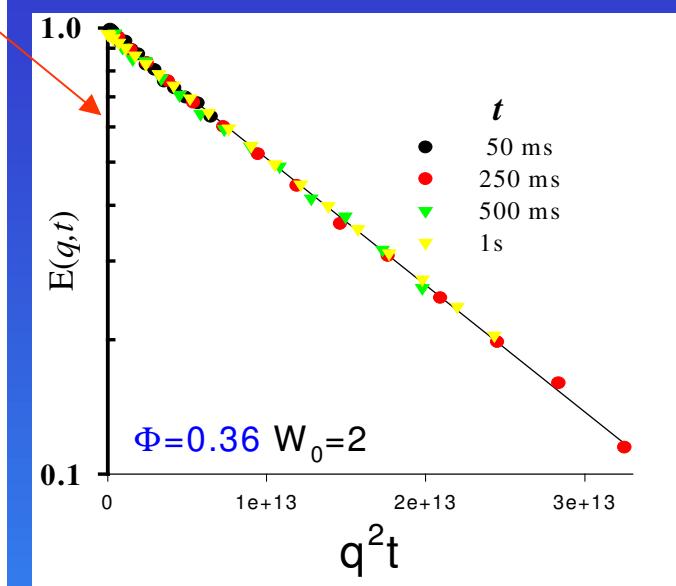
higher concentration: Gaussian diffusion in the whole timescale range



Gels in iC8; $W_0=2$

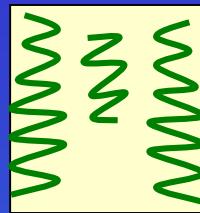
Lecithin mean squared displacement as a function of observation time

More connections



Increasing the W_0 : Gaussian diffusion also at very low volume fraction

Phase separation

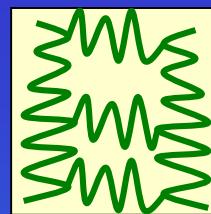


branches



W_0

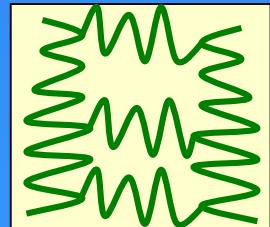
η decreases



more branches



W_0



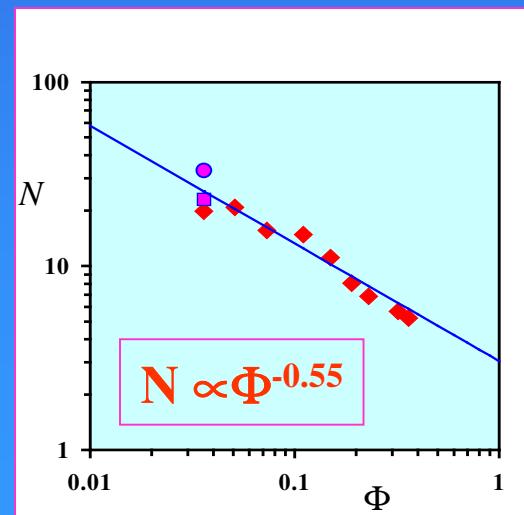
$$D = \frac{D_c}{3N}$$

Theory

Drye & Cates *J. Chem. Phys.* **1992**, 96, 1367

Tlusty & Safran *Phys. Rev. Lett.* **2000**, 84, 1244

Zilman & Safran *Phys. Rev. E* **2002**, 66, art. N. 051107



$$N \propto L_{\text{branch}} \propto \Phi^{-1/2}$$

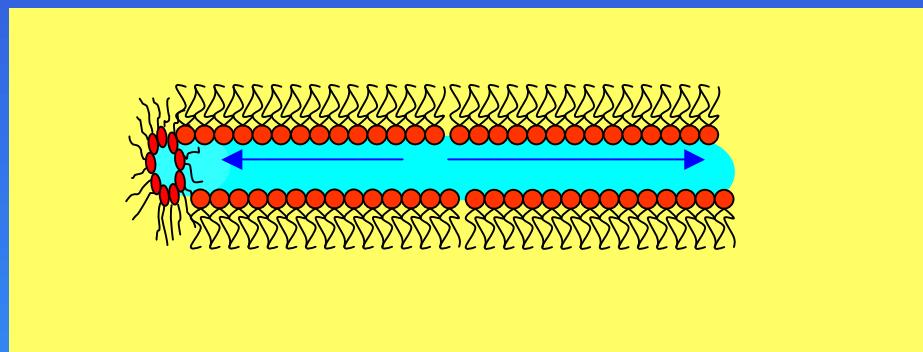
Cyclohexane: water diffusion

For water the residence time is $\sim 1\mu\text{s}$ (Gaussian diffusion)

Water self-diffusion coefficient D_w

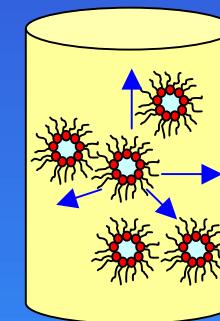
giant slow moving micelles

D_w reflects the water motion inside the micelle



Small fast moving micelles

D_w reflects the overall micelle motion



the motion of water parallel to oriented lipid bilayers was previously studied*

for lamellae D_w increases upon increase in W_0

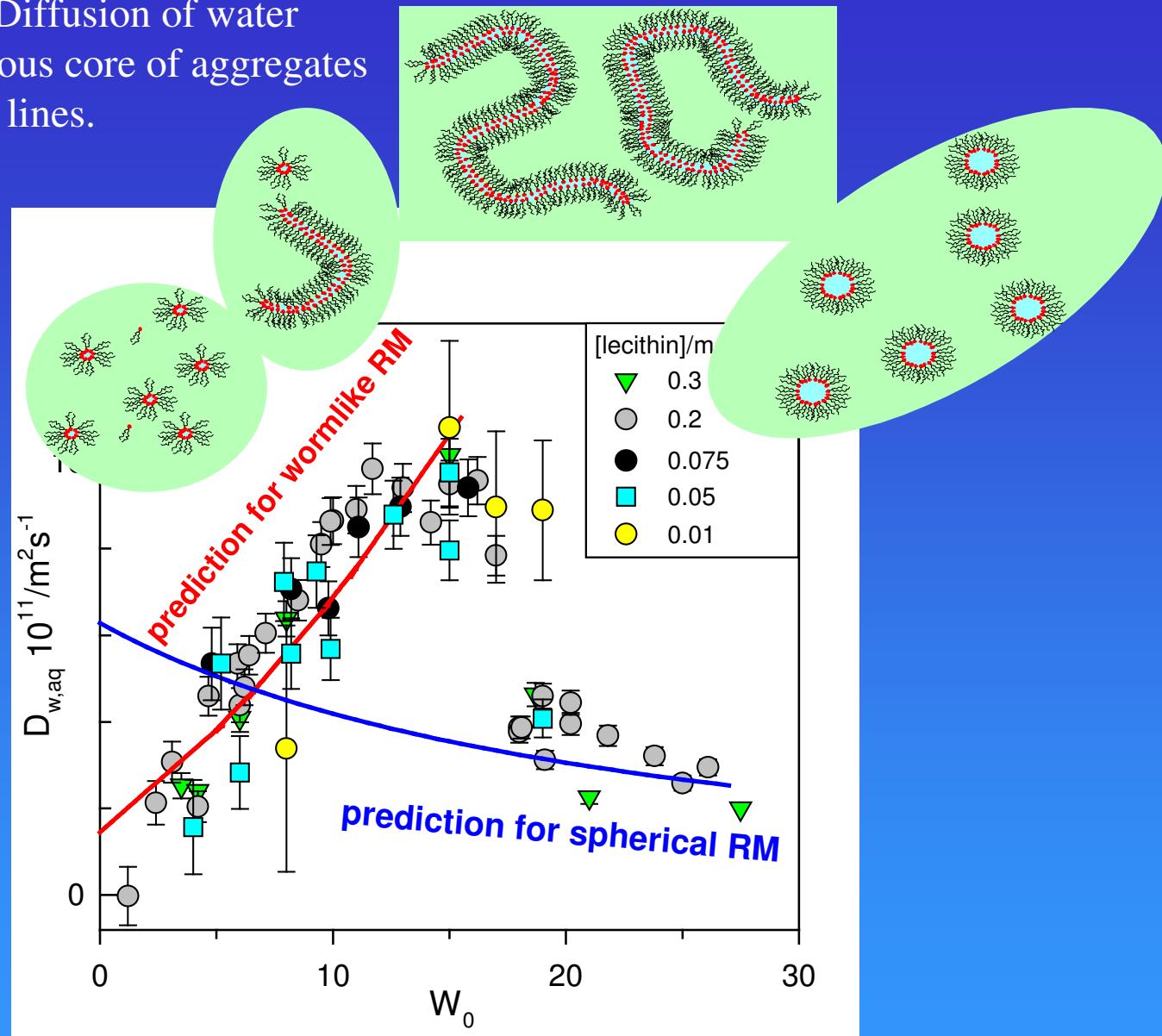
$$D_w(\text{cylinders}) = \frac{1}{3} D_w$$

D_w can be calculated for spherical RM

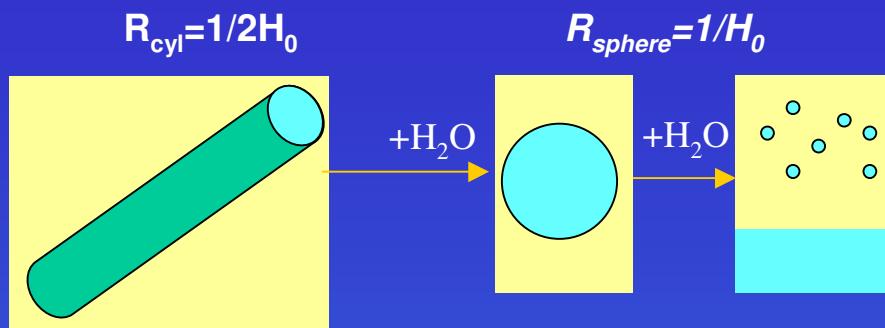
D_w decreases upon increase in W_0

* Wassel *Biophys. J.* 1996, 71, 2724.

Organogels in cC6: Diffusion of water confined in the aqueous core of aggregates along water-dilution lines.



Angelico, Palazzo, Colafemmina, Cirkel, Giustini, Ceglie *J. Phys. Chem. B* 1998, 102, 2883
Angelico, Balinov, Ceglie, Olsson, Palazzo, Soderman *Langmuir*, 1999, 15, 1679



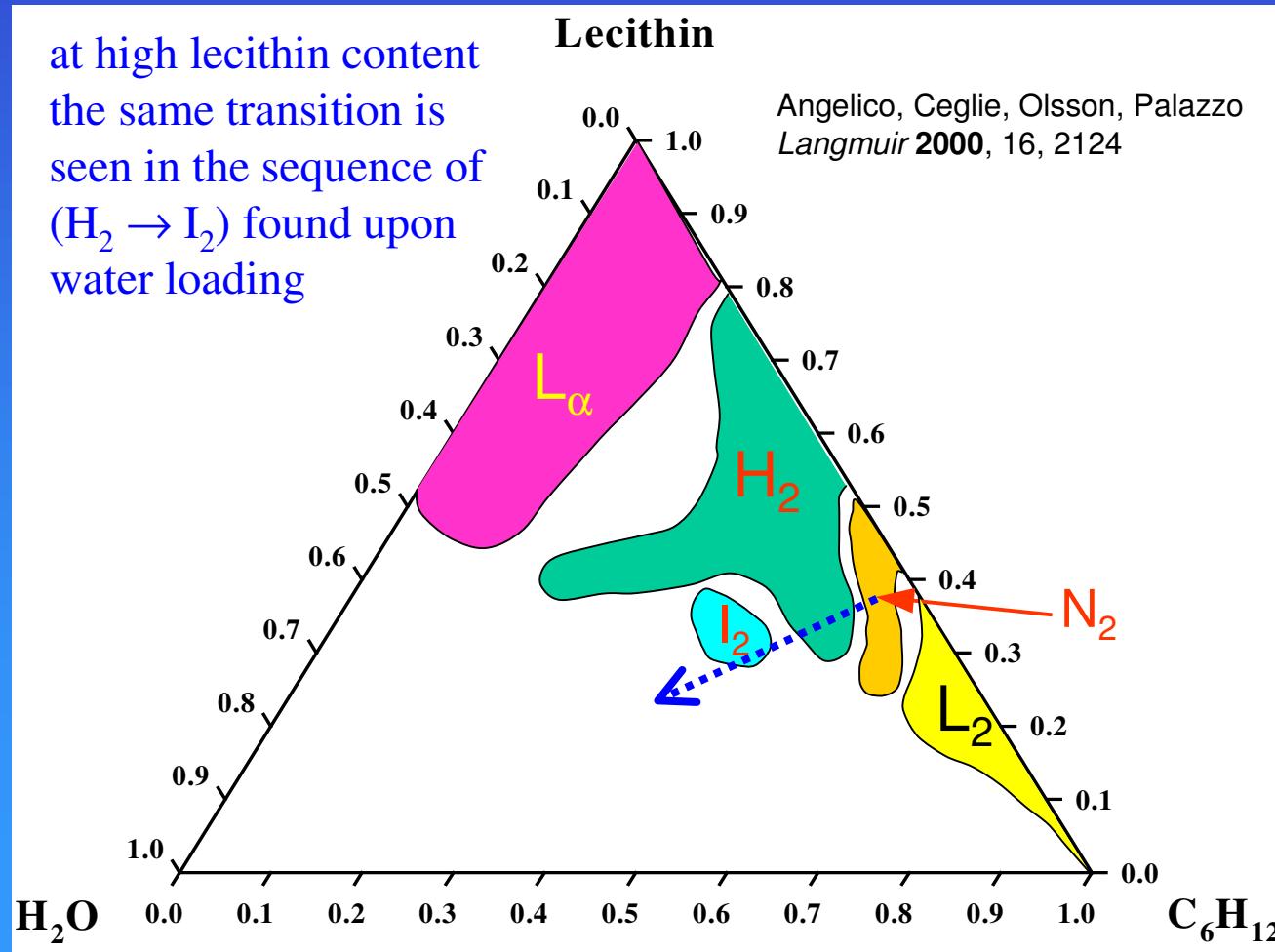
The rod-to-spheres transition accommodates a larger amount of water, leaving constant the mean curvature at the optimal value H_0 .

Upon further water addition the system expels excess water

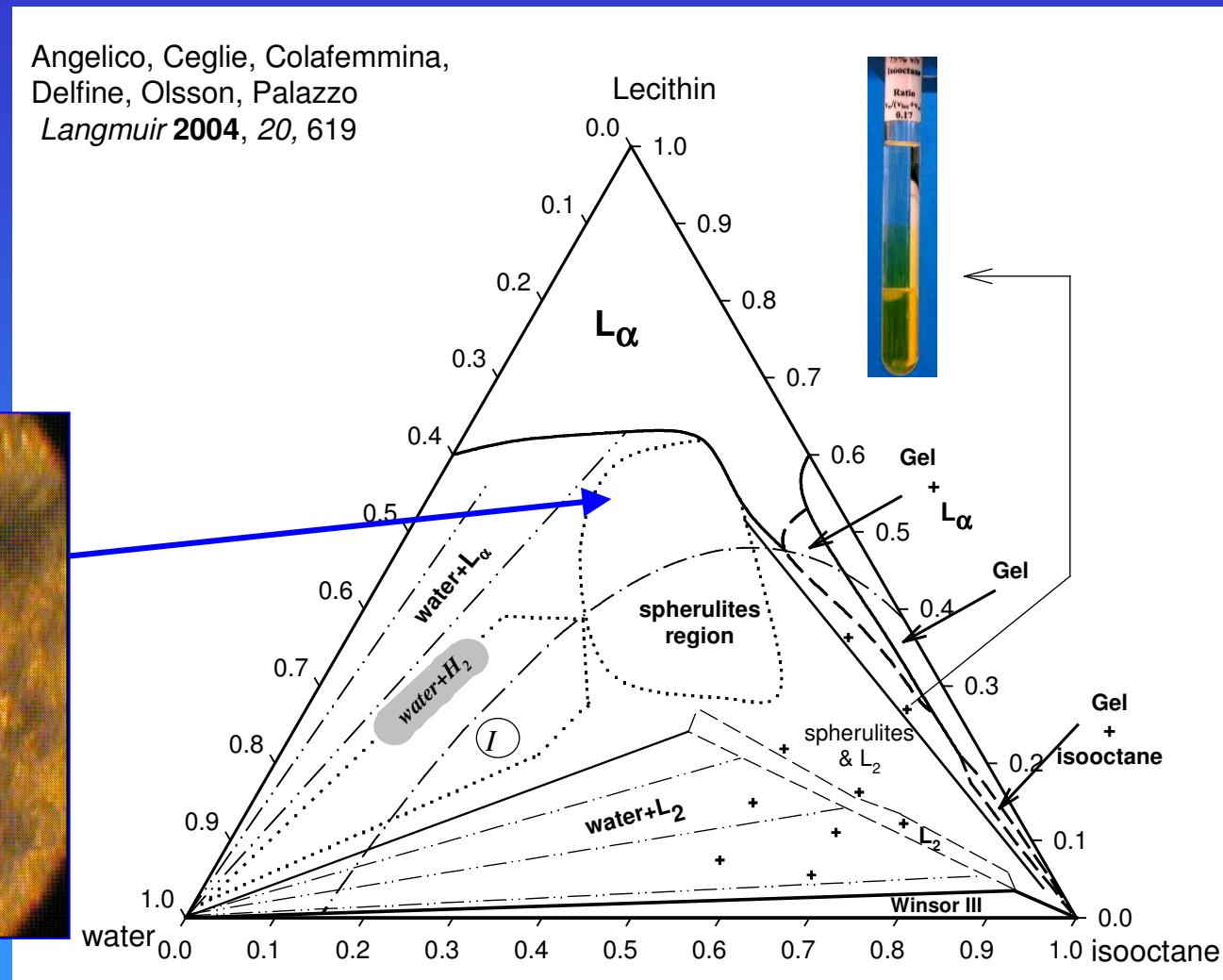
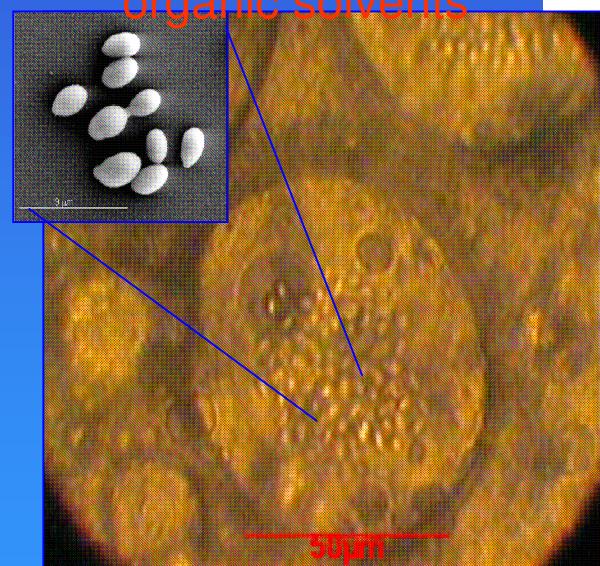
at high lecithin content
the same transition is
seen in the sequence of
($H_2 \rightarrow I_2$) found upon
water loading

Lecithin

Angelico, Ceglie, Olsson, Palazzo
Langmuir **2000**, 16, 2124



For isoctane and linear alkanes the system is governed by *local* curvature energy and large-scale network fluctuations



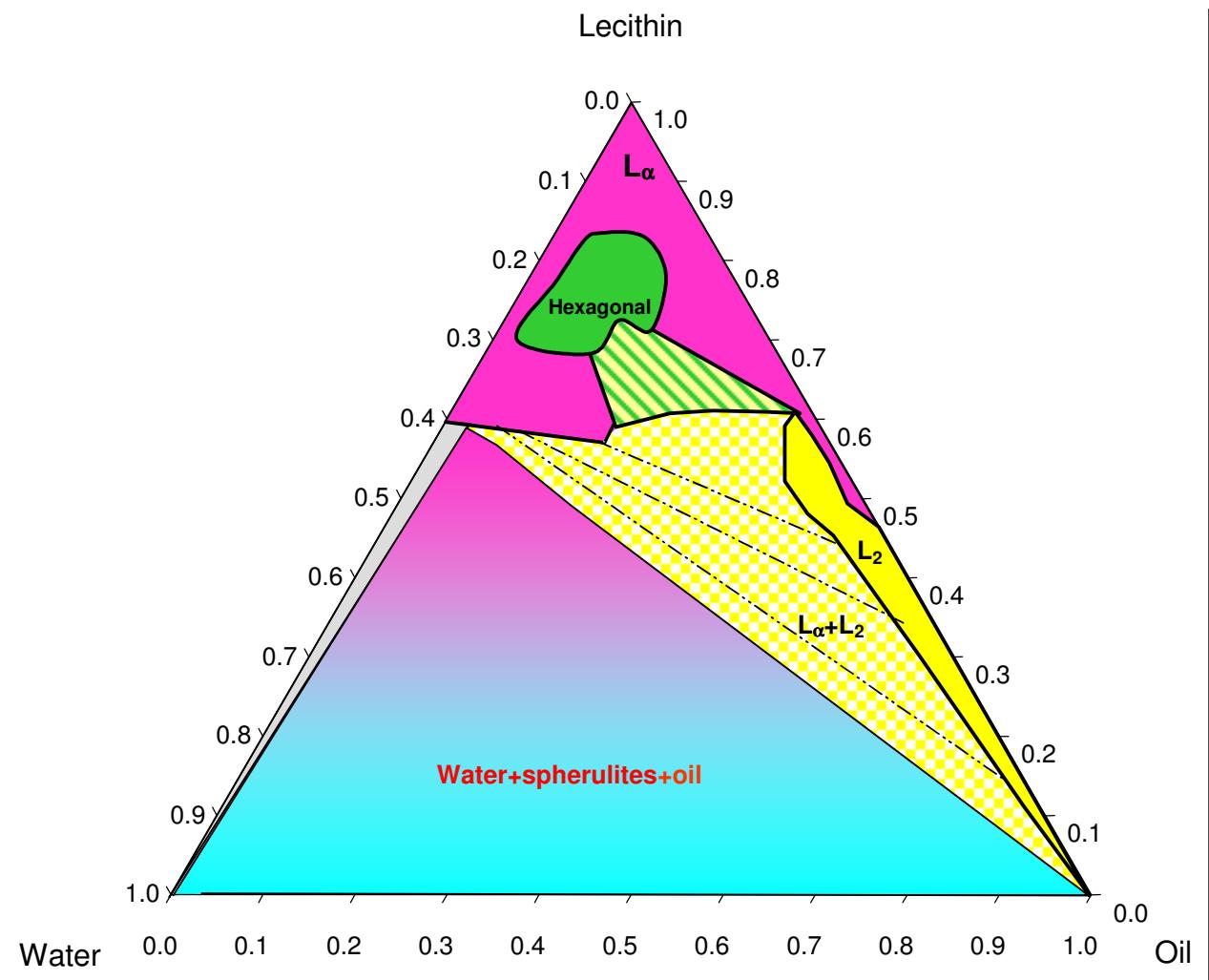
Stefan, Palazzo, Ceglie, Panzavolta, Hochkoepller
Biotech. Bioeng. **2003**, 81, 323.8

Biocompatible fatty acid esters

Use as controlled drug release system

Mackeben, Müller, Müller-Goymann
Colloids Surf. A 2001, 699, 183.

Angelico, Ceglie, Colafemmina, Lopez, Olsson, Palazzo
Langmuir 2005, 21, 140



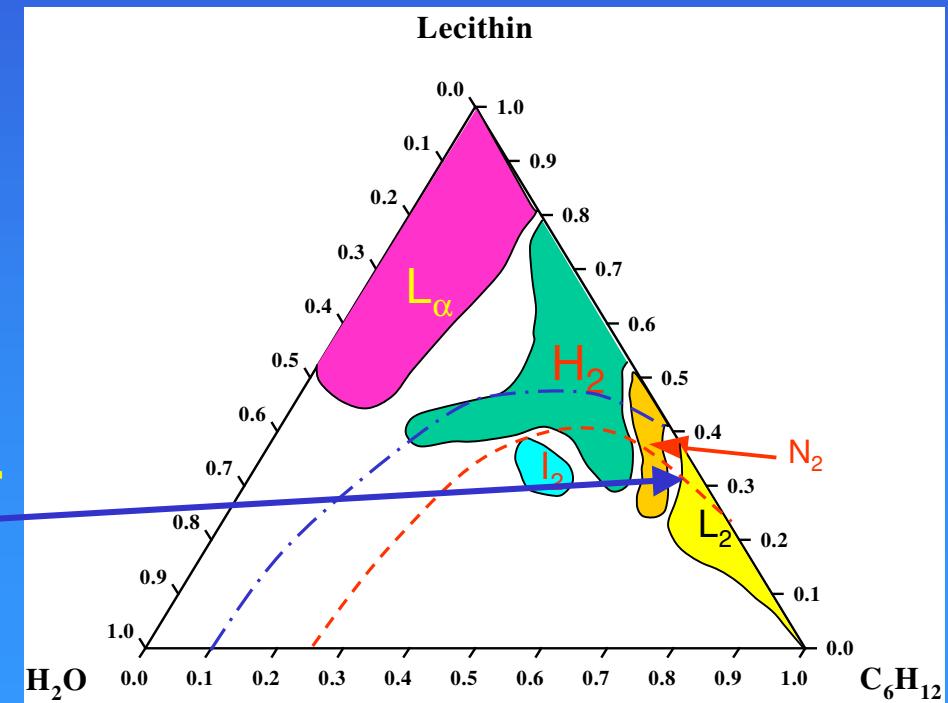
Micellar Dynamics

water residence time of about $1\mu\text{s}$

micellar lifetime and lecithin residence time $> 1 \text{ s}$

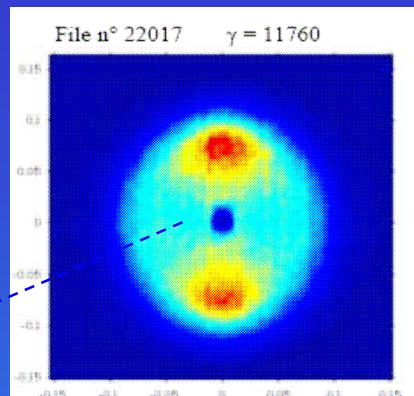
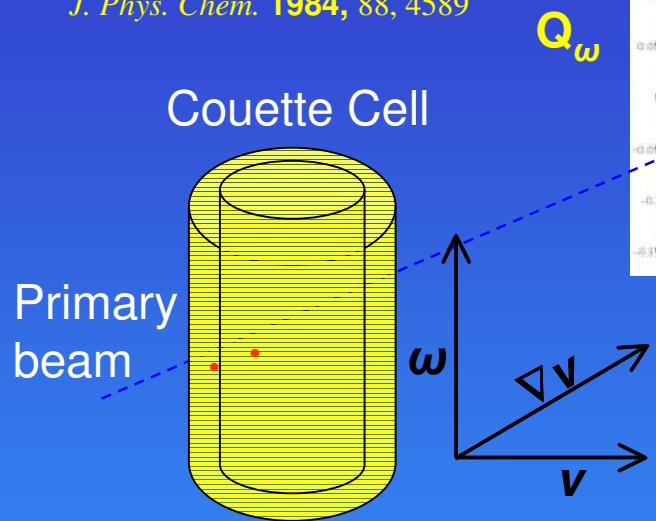
A more precise estimation of upper limit for the micellar lifetime can be obtained in cyclohexane

For isotropic samples close to the boundary of L_2 phase a moderate shear is sufficient to induce the alignment of the micelles into a nematic state.

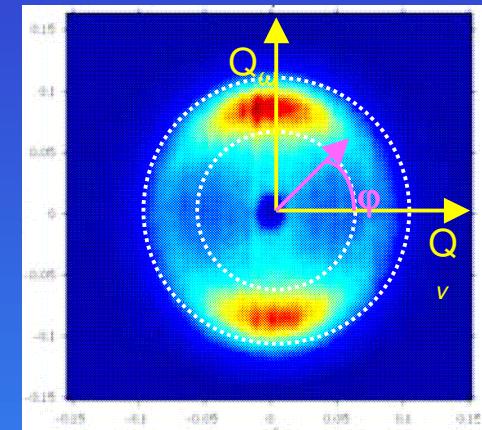


Relaxation of shear-aligned reverse micelles: SANS experiments

Setup:
Hayter, Penfold
J. Phys. Chem. **1984**, 88, 4589



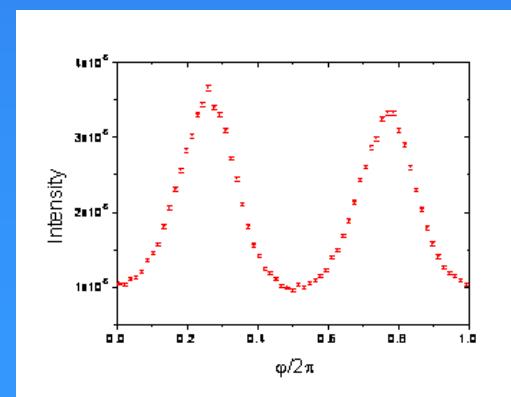
Linear micelles scatter mainly in a direction perpendicular to their long axis.



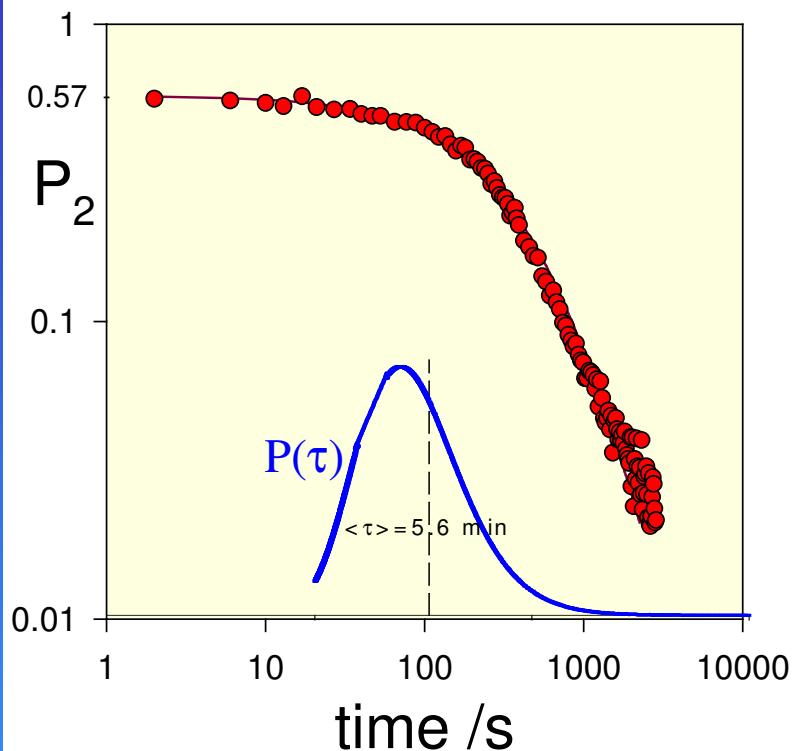
The anisotropy can be analyzed by plotting the scattered intensity as a function of the azimuthal angle, ϕ

from this azimuthal representation, the orientational order can be quantified in terms of the second rank order parameter P_2

Deutsch, *J. Phys. Rev. A* **1991**, 44, 8264

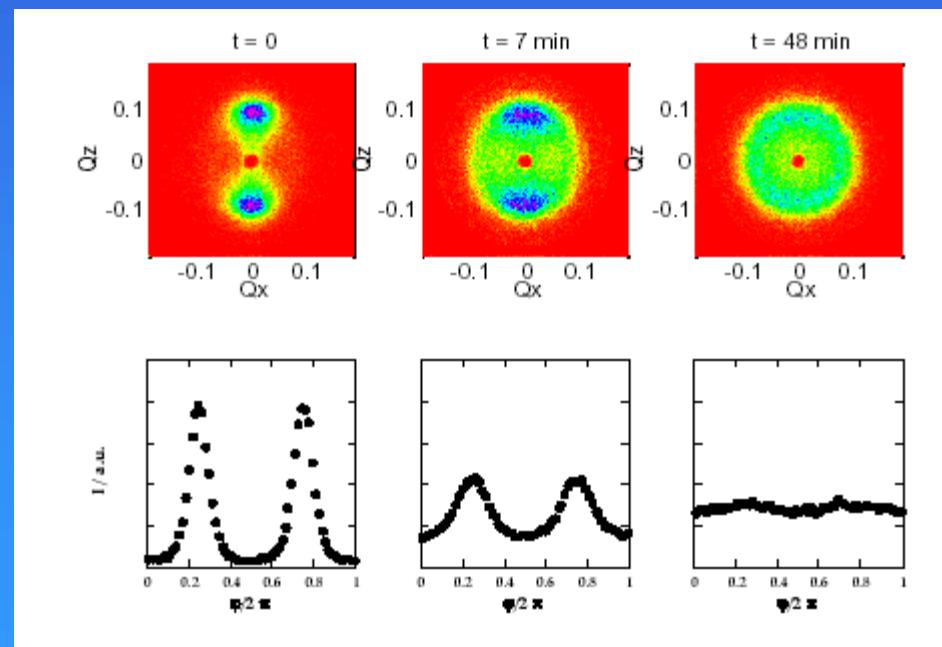


Variation of the second rank order parameter with time after cessation of shear ($dy/dt = 10 \text{ s}^{-1}$). A fit to $P_2 = P_2(0)(1+t/\langle\tau\rangle)^{-n}$ is shown as solid line.



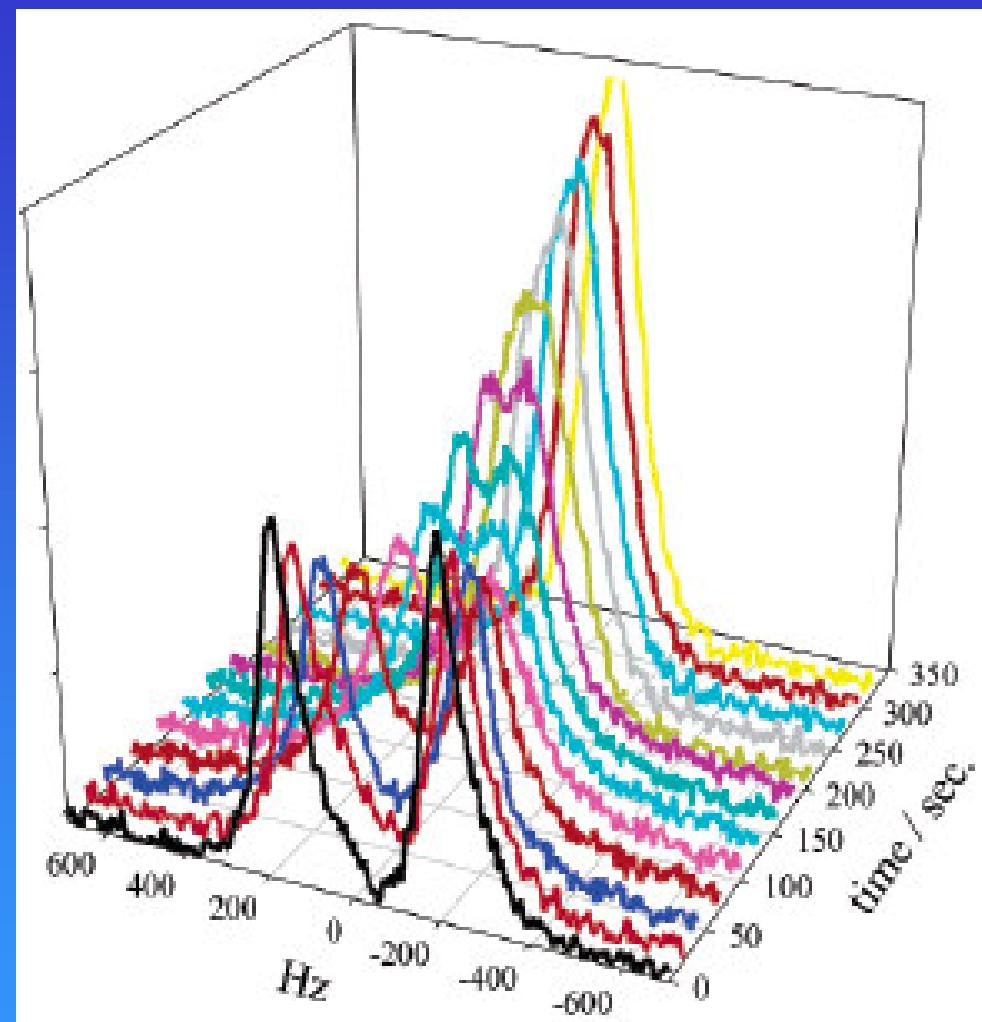
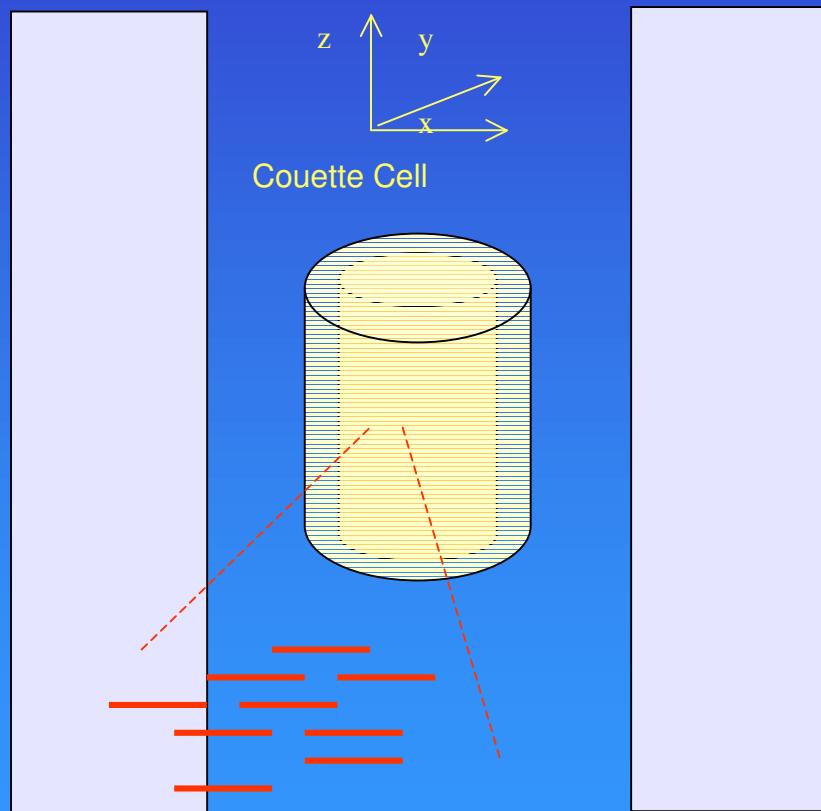
By SANS one can follows the decay of P_2 . The time scale of the order parameter decay is of the order of minutes

The relaxation from N_2 to L_2 is characterized by a gradual decay of orientational order (no nucleation and growth)



Relaxation of shear-aligned reverse micelles: Rheo-NMR experiments

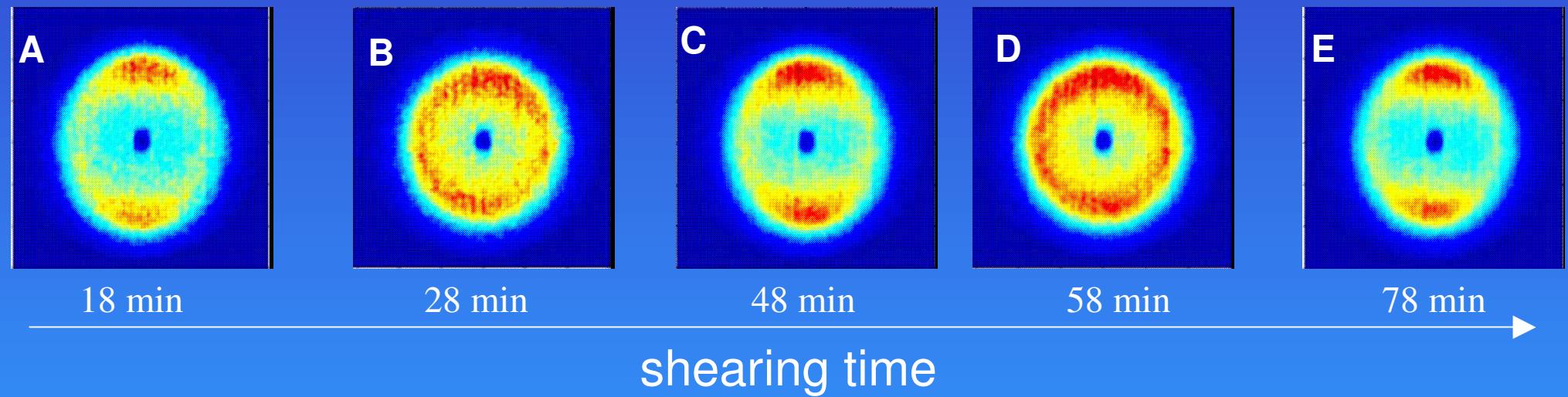
experimental

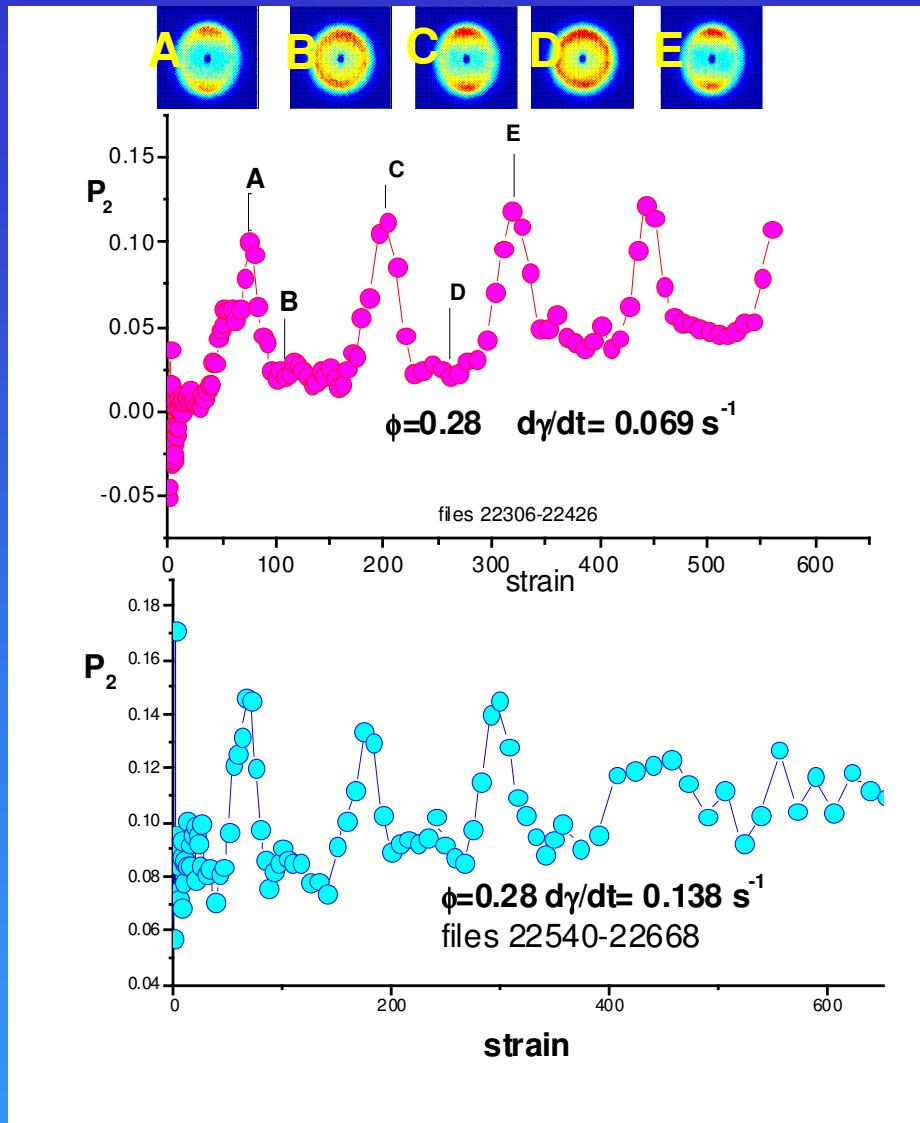


Angelico, Burgemeister, Ceglie, Olsson, Palazzo, Schmidt
J. Phys. Chem. B 2003, 107, 10325

Rheo-SANS start-up experiments at very low shear rate

$\phi=0.28$ $d\gamma/dt = 0.069$ s⁻¹





P_2 exhibits damped oscillations scaling with the strain

Working hypothesis:

oscillations reflect the tumbling dynamics of nematic polydomains during the shear-induced isotropic-nematic transition

Caveat:
more measurements are required

acknowledgements

Ruggero ANGELICO & Andrea CEGLIE – Campobasso, Italy ***grazie!!!***

Kell MORTENSEN - Copenhagen, Denmark ***tak!!!***

Claudia SCHMIDT - Paderborn, Germany ***danke!!!***

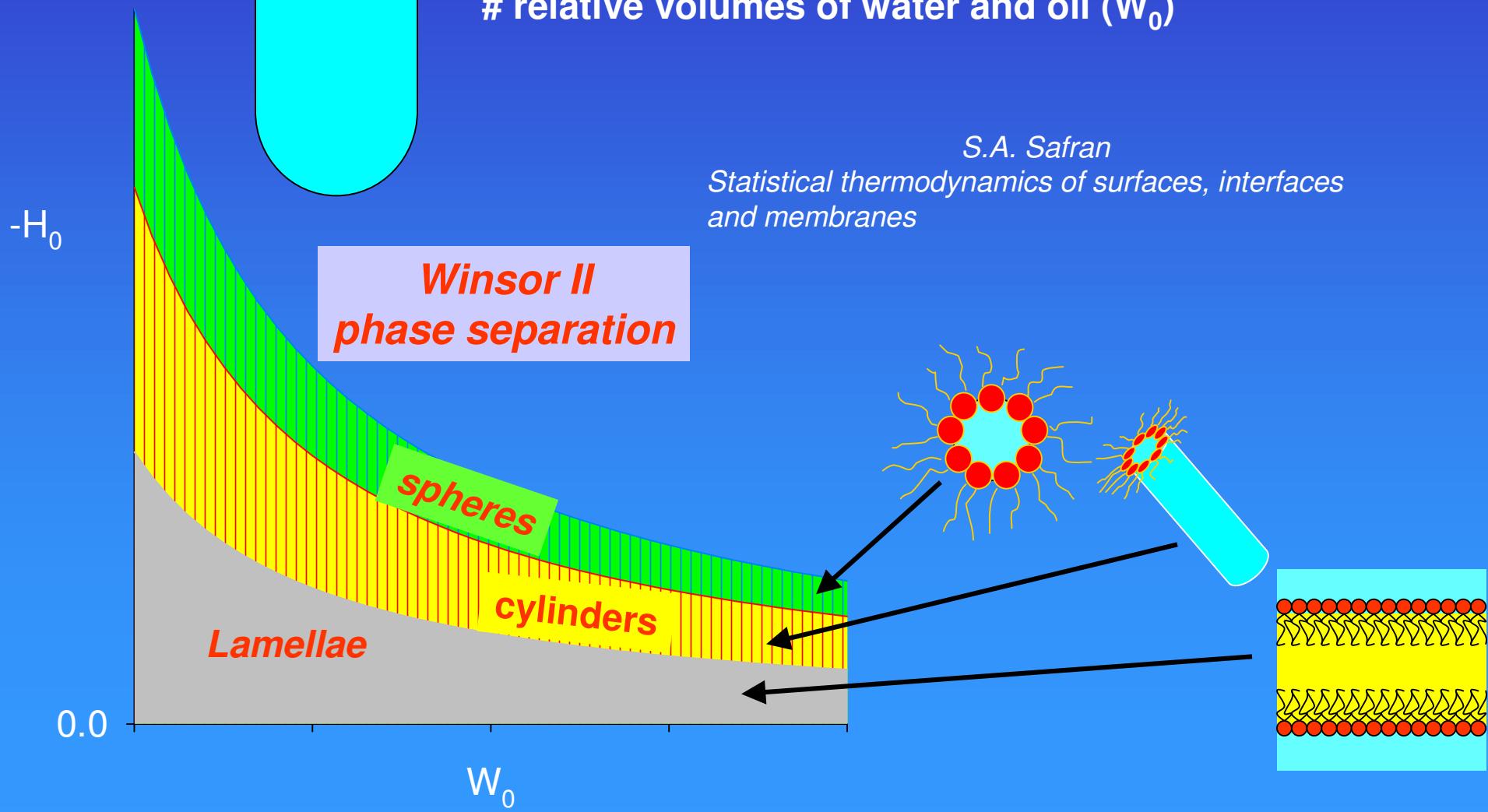
Ulf OLSSON – Lund, Sweden ***tack!!!***

Effect of oil

Aggregates shapes is dictated by:

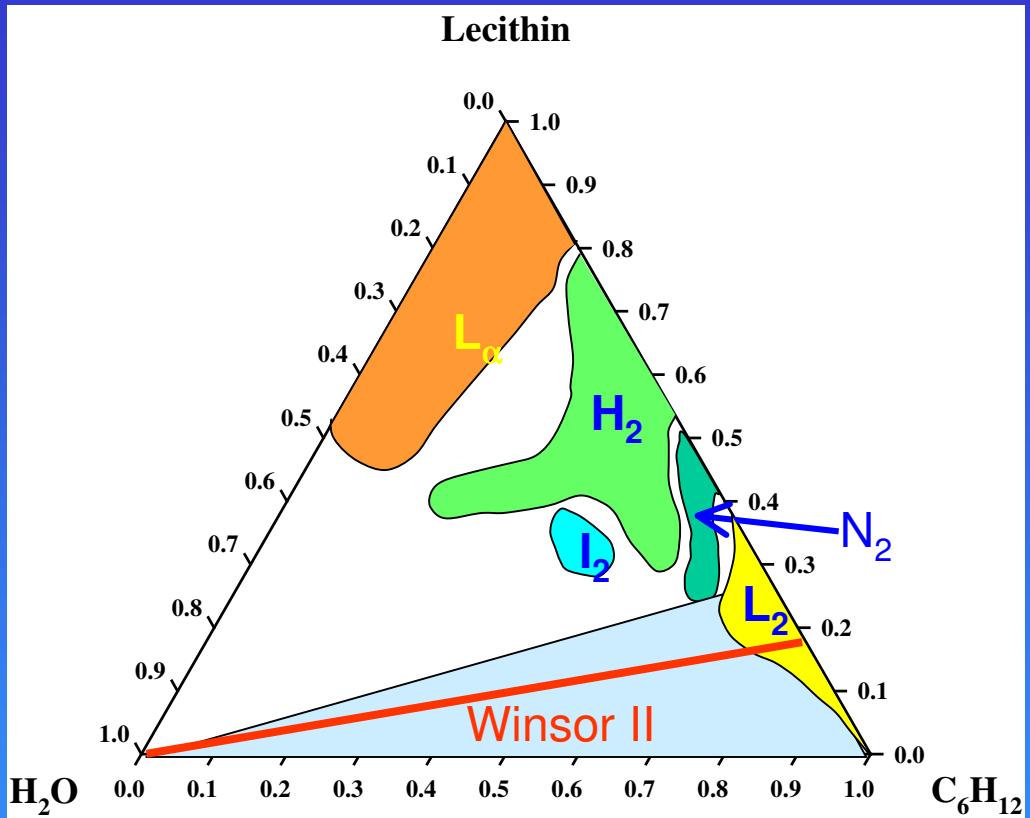
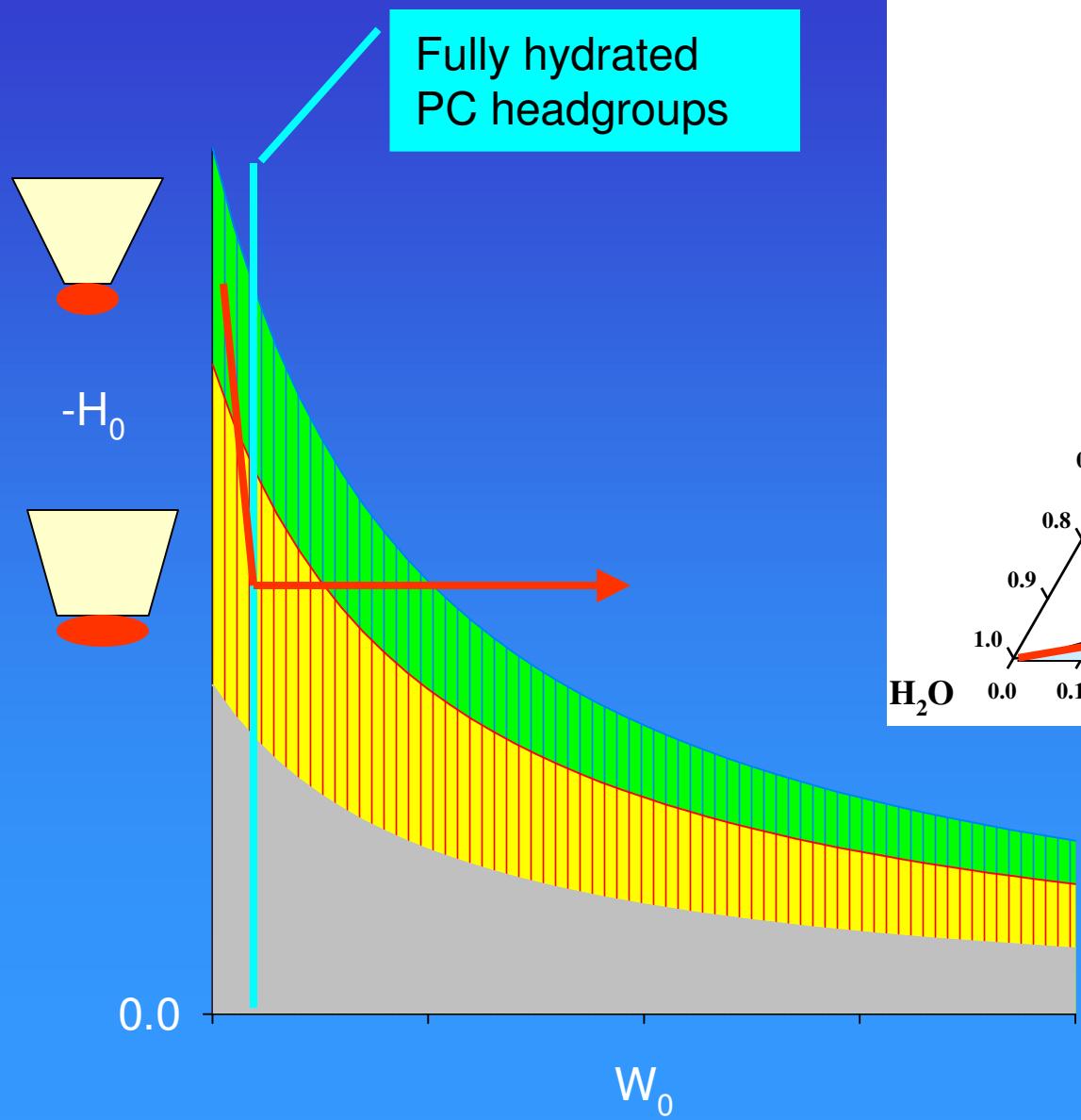
H_0 function of polar head hydration and oil penetration

relative volumes of water and oil (W_0)



Angelico, Ceglie, Olsson, Palazzo
Langmuir 2000, 16, 2124

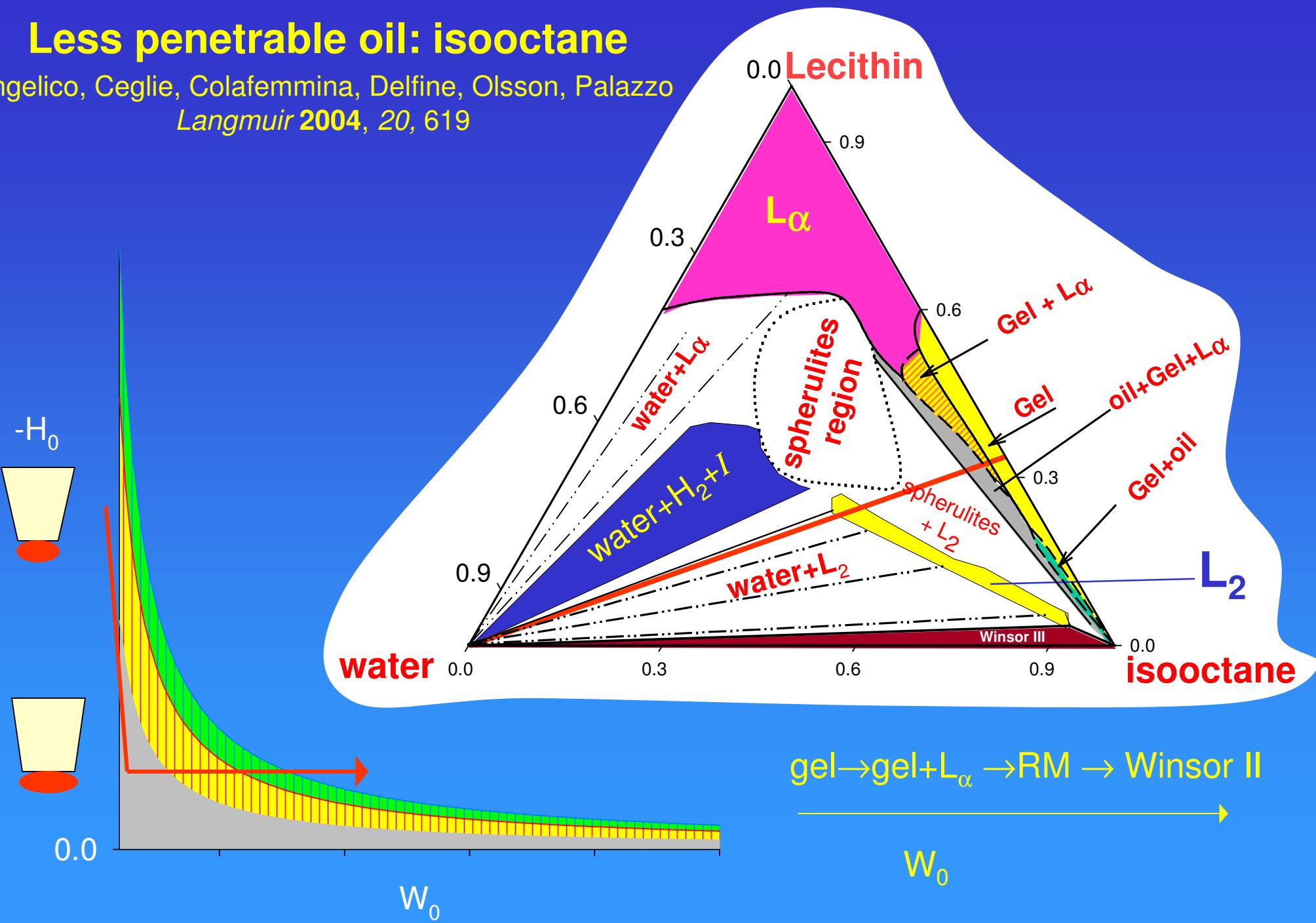
Highly penetrable oil:cyclohexane



Less penetrable oil: isooctane

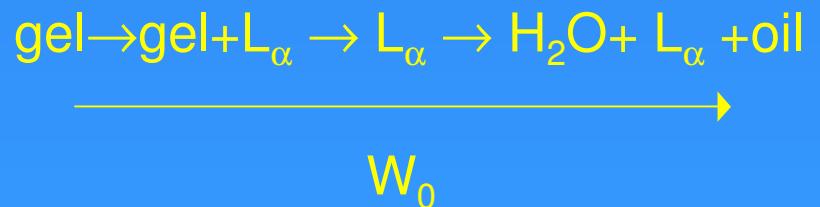
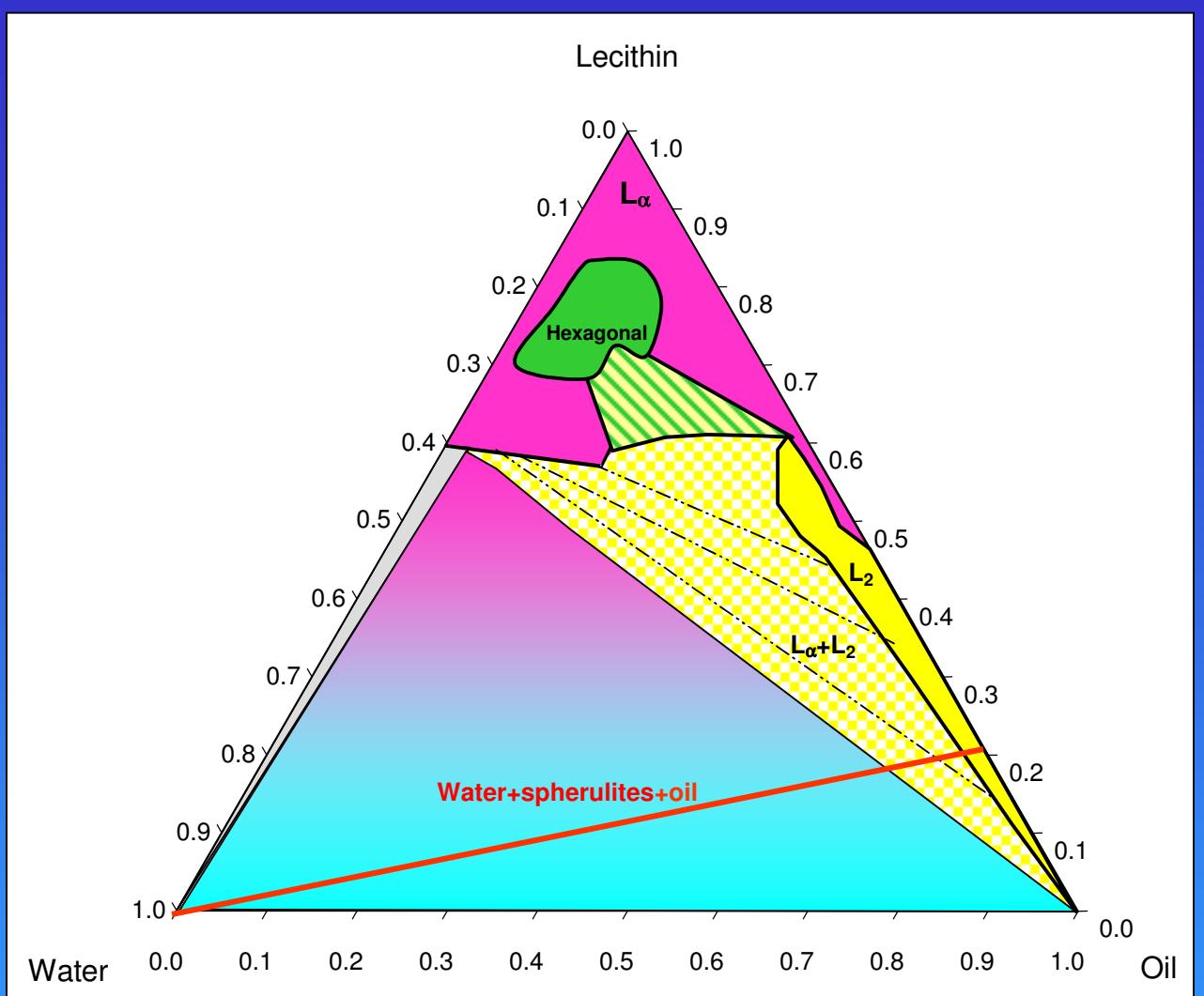
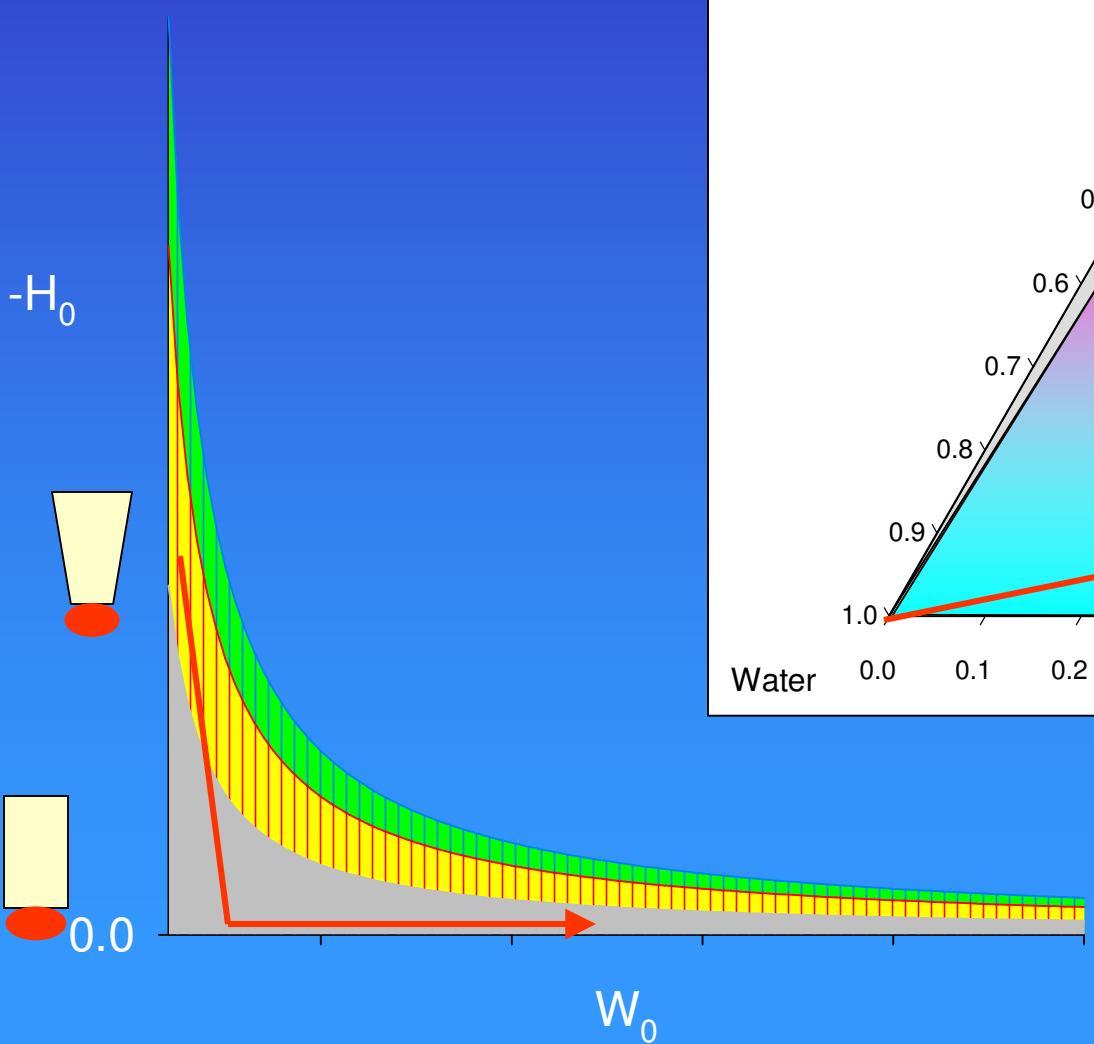
Angelico, Ceglie, Colafemmina, Delfine, Olsson, Palazzo

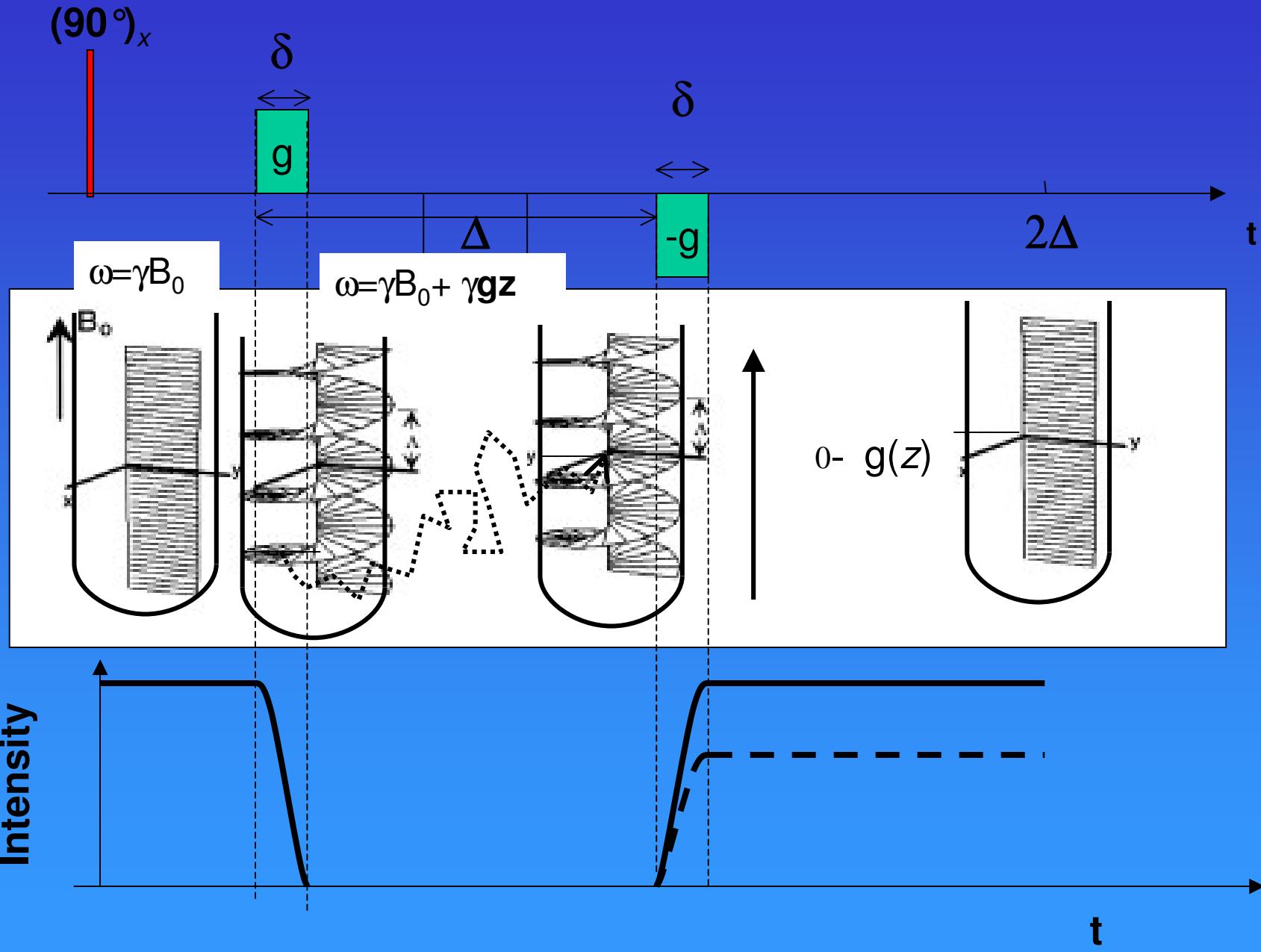
Langmuir 2004, 20, 619



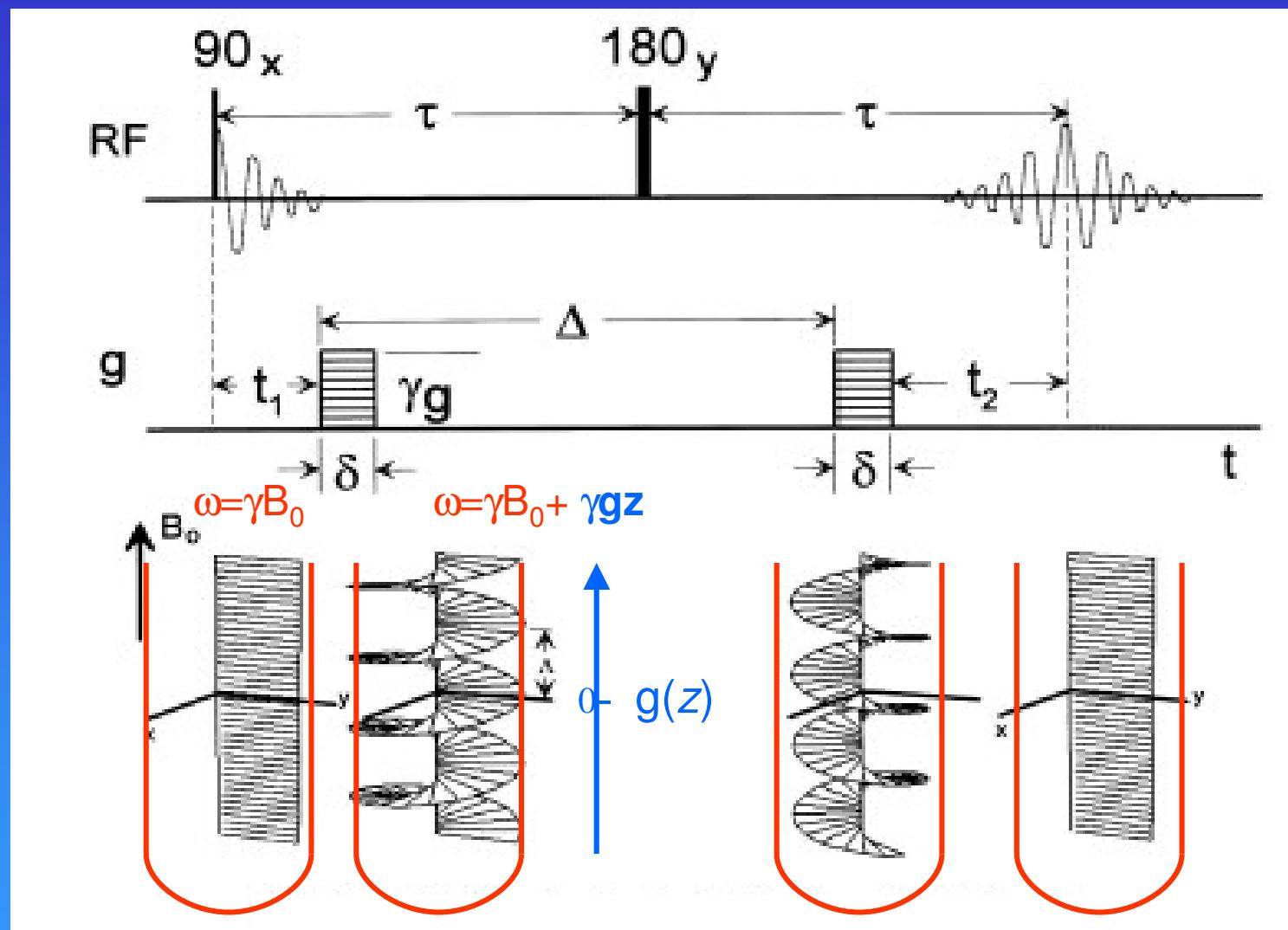
Almost non-penetrable oil: isopropylpalmitate

Angelico, Ceglie, Colafemmina, Lopez, Olsson,
Palazzo
Langmuir 2005

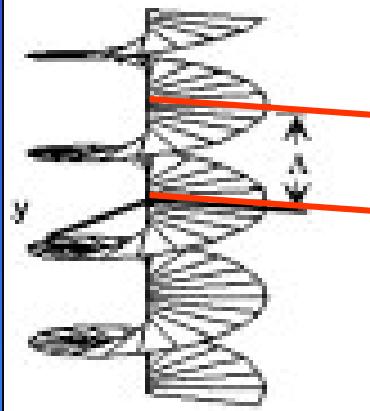




Determinazione di D via PGSE-NMR



$$E(q,t) = \int_{-\infty}^{+\infty} \overline{P}(z,t) e^{iqz} dz$$



For a simple liquid

$$\overline{P}(z,t) = (4\pi Dt)^{-\frac{1}{2}} e^{-\frac{z^2}{4Dt}}$$

$$E(q,t) = e^{-q^2 Dt}$$

$$q = \gamma g \delta$$

$$2\pi/q$$

